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## Magneto-optical Kerr effect in Fe/Au superlattices modulated by integer atomic layers

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### Abstract

We have investigated the magneto-optical Kerr effect (MOKE) in Fe( $n$  ML)/Au( $n$  ML) superlattices where  $n$  is integer and ML represents monatomic layer thickness. The MOKE spectra show prominent structure between 3 and 5 eV. The  $n$  dependence of the prominent structure is discussed in comparison with the case of ultrathin Fe films which show the magneto-optical transitions associated with the quantum-well states. © 1998 Elsevier Science B.V. All rights reserved.

**Keywords:** Magneto-optical effects; Superlattices artificial; Quantum well; Thin films magnetic

Recent development of thin-film preparation techniques has made it possible to fabricate ordered alloys with layered structures artificially. In previous papers [1-3], we reported the artificial fabrication of a L1<sub>0</sub> ordered structure by alternate deposition of (0 0 1)Fe and (0 0 1)Au monatomic layers. The L1<sub>0</sub> ordered structure was successfully obtained although it does not exist in the equilibrium phase diagram. The ordered FeAu alloy film possesses large Fe moment exceeding 2.2  $\mu_B$ , large perpendicular magnetic anisotropy along the  $c$ -axis and novel magneto-optical Kerr spectra.

In this paper, we prepared Fe/Au superlattices modulated by integer atomic layers, i.e., Fe( $n$  ML)/Au( $n$  ML) superlattices where  $n$  is integer and ML represents monatomic layer thickness. The  $n$  dependence of the magneto-optical Kerr spectra is shown and discussed, particularly, compared to the results of ultrathin (0 0 1)Fe films sandwiched with Au [4, 5].

The samples were prepared on polished MgO(1 0 0) substrates by UHV deposition. Fe  $n$  ML and Au  $n$  ML were alternately deposited on a thick Au buffer.  $n$  was varied in the range from 1 to 12. The repetition number,  $N$ , was 100/ $n$ , or the integer closest to 100/ $n$  when 100/ $n$  is

non-integer. The spectra of the Kerr rotation  $\theta_K$  and ellipticity  $\eta_K$  were measured between 1.2 and 6.0 eV by means of polarization modulation technique.

The  $\theta_K$  and  $\eta_K$  spectra for  $n = 1, 2, 3, 4, 6, 8$  and 12 are shown in Fig. 1a and Fig. 1b, respectively. For  $n \geq 6$ , a sharp peak of  $\theta_K$  and a dispersion-type structure of  $\eta_K$  are observed around 2.4 eV, which are caused by the enhancement of magneto-optical signals by plasma resonance of conduction electrons in Au layers [6]. On the other hand, the enhancement is not clearly observed for  $n \leq 4$ , suggesting a considerable change in the electronic structure of Au layers.

Another remarkable feature is a definite dispersion-type structure between approximately 3 and 5 eV in the  $\theta_K$  spectra for  $1 \leq n \leq 8$ . Correspondingly, the  $\eta_K$  spectra show the peak around the dispersion center of  $\theta_K$ . These prominent structures shift to higher energies as  $n$  increases from 1 to 6. Fig. 2 shows the spectra of the absolute values of  $\omega\sigma_{xy}$ , the off-diagonal element of the complex conductivity tensor multiplied by angular frequency. The  $\omega\sigma_{xy}$  spectra were obtained from the experimental data of  $\theta_K$ ,  $\eta_K$  and the diagonal dielectric constant,  $\epsilon_{xx}$ , evaluated by the Kramers-Kronig analysis of the reflectivity spectra. The arrows in Fig. 2 indicate the peaks corresponding to the prominent structures observed in the  $\theta_K$  and  $\eta_K$  spectra. The Fe-layer thickness  $d_{Fe}$  and/or  $n$  dependence of the position of the peak in the

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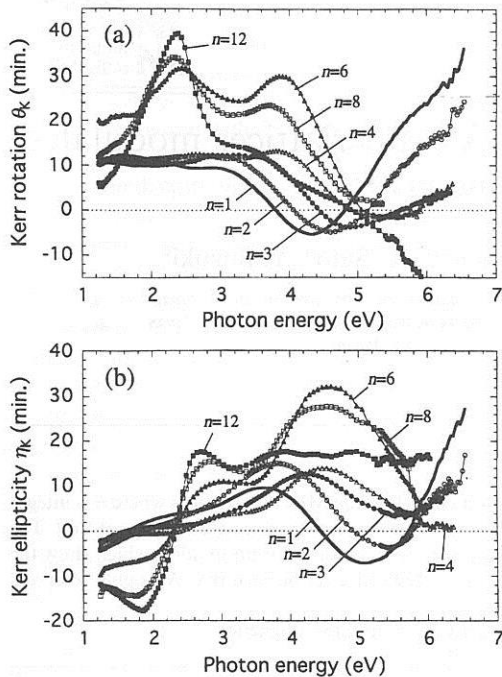


Fig. 1. (a) Magneto-optical Kerr rotation  $\theta_K$  and (b) ellipticity  $\eta_K$  spectra for Fe( $n$  ML)/Au( $n$  ML) superlattices.

$\omega\sigma_{xy}$  spectra is shown in Fig. 3 (○). The position of the peak shifts from approximately 3.2 to 4.4 eV as  $n$  increases from 1 to 4. For  $n \geq 4$ , however, it shows no considerable change. This result is compared to the case of ultrathin Fe films. Suzuki et al. [4] reported the appearance of new structures in the magneto-optical Kerr spectra of ultrathin (001)Fe films sandwiched with Au, and they associated the new structures to transitions from the  $\Delta_5$  band to quantum well states in the  $\Delta_1$  band of Fe films [5]. The transition energy obtained from the  $\omega\sigma_{xy}$  spectra of the ultrathin Fe films are also shown in Fig. 3 (●). The  $d_{Fe}$  dependence of the  $\omega\sigma_{xy}$  peak in Fe/Au superlattices has similar tendency to that of the transition energy associated with the quantum-well states in ultrathin Fe films. This suggests the formation of quantum-well states within Fe layers even for Fe/Au superlattices, which causes the peak of  $\omega\sigma_{xy}$ . Quantitatively, however, the transition energies in Fe/Au superlattices are somewhat higher than those in ultrathin Fe films. Particularly, considerable discrepancy is seen for  $n < 3$ . We consider that for  $n < 3$  the electrons cannot be confined completely within Fe layers because Au layers are also very thin, and the quantum-well states are spread into minibands, leading to the discrepancy between the transition energies in ultrathin Fe films and in Fe/Au superlattices.

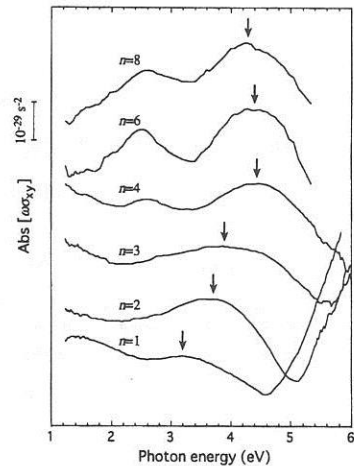


Fig. 2. Absolute values of  $\omega\sigma_{xy}$ , the off-diagonal element of the complex conductivity tensor multiplied by angular frequency, for Fe( $n$  ML)/Au( $n$  ML) superlattices as a function of photon energy. The arrows indicate the peaks corresponding to the prominent structures observed in the  $\theta_K$  and  $\eta_K$  spectra (Fig. 1).

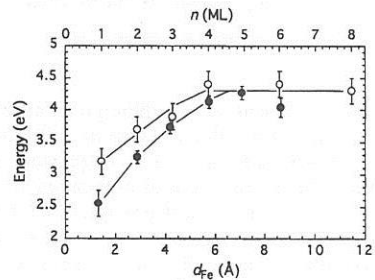


Fig. 3. Fe layer thickness  $d_{Fe}$  and/or  $n$  dependence of the position of the peak in the  $\omega\sigma_{xy}$  spectra for Fe( $n$  ML)/Au( $n$  ML) superlattices (○). For comparison, the magneto-optical transition energies associated with quantum-well states in ultrathin (001)Fe films sandwiched with Au[6] are shown (●).

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